

Find \vec{B} by integration: $\vec{B} = \frac{\mu_0}{4\pi} \int \frac{I d\vec{\ell}' \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$

Note: $I d\vec{\ell}$, $\vec{K} dA$, $\vec{J} dV$ depending ...

\vec{r} : location at which \vec{B} is being found

\vec{r}' : location of a source - i.e. current

$d\vec{\ell}$, \vec{K} , \vec{J} : direction current is going

Long straight wire: $\vec{r} = (x, 0, 0)$ $\vec{r}' = (0, 0, z')$
 $d\vec{\ell} = (0, 0, dz')$

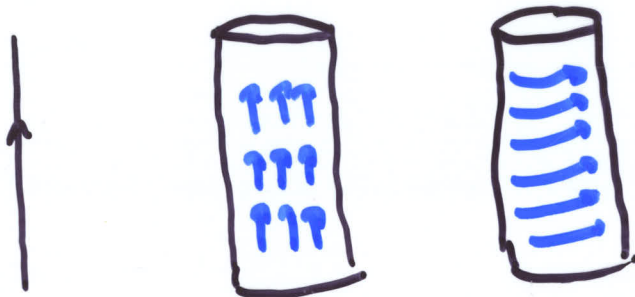
$\vec{B} = \frac{\mu_0 I}{2\pi r} \hat{\phi}$ in cylinders from z axis where wire is

Cylindrical Shell: $\vec{K} dA = (0, 0, K) R d\theta dz'$
 $\vec{r}' = (R \cos\theta, R \sin\theta, z') \sim I/2\pi R$

outside: \vec{B} as above; inside $\vec{B} = 0$

Solenoid: $\vec{K} dA = K(-\sin\theta, \cos\theta, 0) R d\theta dz'$

outside: $\vec{B} = 0$; inside $\vec{B} = \mu_0 K \hat{k} = \frac{\mu_0 N}{l} I \hat{k}$



Note use of odd symmetry to show integral = 0

Dwight 200.03

$$\int \frac{dx}{r^3} = \frac{1}{a^2} \frac{x}{r}$$

$$r = \sqrt{a^2 + x^2}$$

859.124

$$\int_{-\pi}^{\pi} \frac{a - b \cos x}{a^2 - 2ab \cos x + b^2} dx = \begin{cases} \frac{2\pi}{a} \\ 0 \end{cases}$$

if $\begin{cases} a > b \\ a < b \end{cases}$

$$\vec{E} = -\vec{\nabla} \phi \leftarrow +C$$

$$\vec{B} = \vec{\nabla} \times \vec{A} \leftarrow -\vec{\nabla} \phi$$

↑
vector pot

$$\vec{B} = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(r') \times (\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} dV'$$

-∇ $\frac{1}{|\vec{r} - \vec{r}'|}$

p20 1-1-9

$$\nabla \times (\phi \vec{F}) = \nabla \phi \times \vec{F} + \phi \nabla \times \vec{F}$$

$$= \frac{\mu_0}{4\pi} \int \nabla \times \left(\frac{J(r')}{|\vec{r} - \vec{r}'|} \right) - \phi \frac{\nabla \times J(r')}{0}$$

$$\vec{B} = \frac{\mu_0}{4\pi} \nabla \times \int \frac{\vec{J}(r')}{|\vec{r} - \vec{r}'|} dV'$$

(\vec{A} Gauge Transform)
↑
Gauge Force

$$\phi = \frac{1}{4\pi\epsilon_0} \int \frac{\rho dV}{|\vec{r} - \vec{r}'|}$$

$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(r')}{|\vec{r} - \vec{r}'|} dV'$$

$$\nabla^2 \phi = -\frac{\rho}{\epsilon_0} \quad \dots \quad \nabla^2 \frac{1}{|\mathbf{r}-\mathbf{r}'|} = -4\pi \delta(\mathbf{r}-\mathbf{r}')$$

$$\nabla^2 A_\gamma = \frac{\mu_0}{4\pi} \int J_\gamma(\mathbf{r}') \underbrace{\nabla^2 \frac{1}{|\mathbf{r}-\mathbf{r}'|}}_{-4\pi \delta} dV'$$

$$\underbrace{\hspace{10em}}_{-4\pi J_\gamma(\mathbf{r})}$$

$$= -\mu_0 J_\gamma$$

$$\nabla^2 \vec{A} = -\mu_0 \vec{J}$$

$$\vec{A} \times (\vec{B} \times \vec{C}) = \vec{B} \cdot \vec{A} \vec{C} - \vec{C} \cdot \vec{A} \vec{B}$$

$$\nabla \times (\nabla \times \vec{C}) = \nabla^2 \vec{C} - \nabla(\nabla \cdot \vec{C})$$

$$\nabla^2 \vec{C} = \nabla \times (\nabla \times \vec{C}) + \nabla(\nabla \cdot \vec{C})$$

$$\nabla \cdot \vec{A} = 0$$

$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} dV'$$

$$\nabla \cdot \mathbf{A} = \frac{\mu_0}{4\pi} \int \nabla \cdot \left(\frac{\mathbf{J}(r')}{|r-r'|} \right) dV'$$

$$\nabla \cdot (\phi \vec{F}) = \underbrace{\nabla \cdot \phi \vec{F}} + \phi \cdot \underbrace{\nabla \cdot \vec{F}}$$

$$= \frac{\mu_0}{4\pi} \int \left(\nabla \frac{1}{|r-r'|} \right) \cdot \mathbf{J}(r') dV'$$

$\uparrow -\nabla' \frac{1}{|r-r'|}$

$$= \frac{-\mu_0}{4\pi} \int \underbrace{\nabla' \frac{1}{|r-r'|} \cdot \mathbf{J}(r')} dV'$$

$$\nabla' \cdot \left(\frac{1}{|r-r'|} \mathbf{J}(r') \right) = \frac{1}{|r-r'|} \underbrace{\nabla' \cdot \mathbf{J}(r')} = 0$$

$$= \frac{-\mu_0}{4\pi} \oint \frac{1}{|r-r'|} \mathbf{J}(r') \cdot \hat{n} dA'$$



$$\begin{cases} \nabla \cdot \mathbf{A} = 0 \\ \nabla^2 \vec{A} = -\mu_0 \mathbf{J} \end{cases}$$

$$\begin{aligned} \nabla \times \mathbf{B} &= \nabla \times (\nabla \times \mathbf{A}) = \underbrace{\nabla} (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A} \\ &= \mu_0 \mathbf{J} \end{aligned}$$

$$\nabla \times (\nabla f) = 0$$

$$\nabla \cdot (\nabla \times F) = 0$$

$\nabla \cdot E = \rho / \epsilon_0$
 $\nabla \times E = 0$ *induct*
 $\hookrightarrow E = -\nabla \phi$

$$\begin{pmatrix} \partial_x & \partial_y & \partial_z \\ \partial_x & \partial_y & \partial_z \\ F_x & F_y & F_z \end{pmatrix} = 0$$

$\nabla \cdot B = 0$
 $\nabla \times B = \mu_0 J$

"proof"

! Coulomb
 ! Biot-Savart X
 static

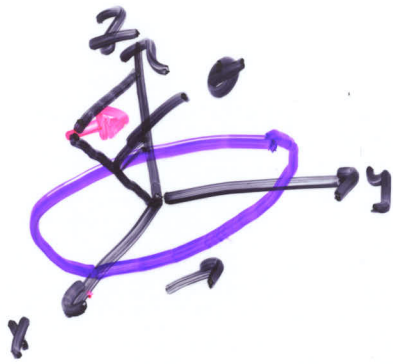
Maxwell's Eq

$\oint E \cdot n dA = \frac{Q_{enc}}{\epsilon_0}$
 $\oint B \cdot dl = \mu_0 I_{enc}$
 Ampere Law



$$\int J \cdot n dA = I = \frac{1}{\mu_0} \int \nabla \times B \cdot dA = \frac{1}{\mu_0} \oint B \cdot dl$$

\vec{A} is not a friend.



$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(r')}{|r-r'|} dV'$$

$$\vec{J} dV \rightarrow \vec{K} dA \rightarrow I d\vec{\ell}$$

$$\vec{r} = (x, 0, z)$$

$$(-\sin\phi, \cos\phi, 0)$$

$$\vec{r}' = (R \cos\phi, R \sin\phi, 0)$$

$$\vec{r}-\vec{r}' = (x-R\cos\phi, -R\sin\phi, z)$$

$$|\vec{r}-\vec{r}'| = \sqrt{\underbrace{x^2}_{\uparrow} + \underbrace{R^2}_{\leftarrow} + \underbrace{z^2}_{\rightarrow} - 2xR\cos\phi}$$

$$\vec{A} = \frac{I\mu_0}{4\pi} \int_{-\pi}^{\pi} \frac{(-\sin\phi, \cos\phi, 0) R d\phi}{|\vec{r}-\vec{r}'|}$$

$$A_\phi = \frac{I\mu_0 R}{4\pi} \int_{-\pi}^{\pi} \frac{\cos\phi d\phi}{\sqrt{r^2 + R^2 - 2xR\cos\phi}}$$

$$K \quad E \quad \frac{2rR\sin\theta \cos\phi}{r^2 + R^2}$$

$$[1-x]^{-1/2} = 1 + \frac{1}{2}x + \dots + \frac{1}{n!} \left(\frac{1}{2}\right)_n x^n + \dots$$

$$A_{\phi} = \frac{I \mu_0 R}{4\pi (r^2 + R^2)^{1/2}} \int_{-\pi}^{\pi} \left[\frac{\cos \phi}{1 + \frac{r^2}{R^2}} \right] d\phi$$

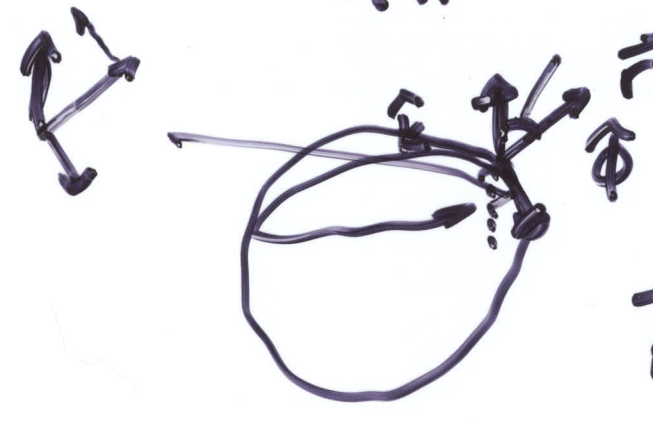
$$\int \cos^2 \phi = \frac{1}{2} \cdot 2\pi \quad \frac{R \sin \theta}{r^2 + R^2} \cos \phi$$

$$= \frac{I \mu_0 R}{4\pi (r^2 + R^2)^{3/2}} \pi \rightarrow \pi R^2 I$$

$\vec{p} = \pm I$
 $q d$
 $m = AI$
 magnetic dipole moment

$$= \frac{\mu_0 m}{4\pi} \frac{r \sin \theta}{(r^2 + R^2)^{3/2}}$$

$$= \frac{\mu_0 m}{4\pi} \frac{\sin \theta}{r^2} \left| \frac{\mu_0 m}{4\pi} \frac{r \sin \theta}{R^3} \right.$$



$$\vec{A} = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \vec{r}}{r^3}$$

$$\vec{r} = A \hat{r} + B \hat{\theta} \quad \vec{r} \cdot \hat{\theta} = -\sin \theta$$