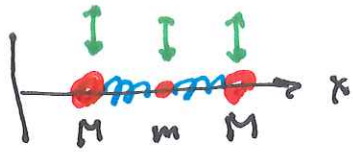


More complex normal mode problems - add wiggle to stretch



add y motion to CO₂ problem

need a restoring force related to

θ angle with an equilibrium

point at $\theta = 180^\circ$ & larger PE as move away from straight.



Cross product of CO bonds provides something

0 when straight and a difference (which we can square) when bent.

$$\vec{r}_1 - \vec{r}_2 + \vec{r}_3 - \vec{r}_2 = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ x_1 - x_2 - l & y_1 - y_2 & 0 \\ x_3 - x_2 + l & y_3 - y_2 & 0 \end{vmatrix} = (x_1 - x_2 - l)(y_3 - y_2) - (x_3 - x_2 + l)(y_1 - y_2)$$

neglect $x \cdot y$

$$= -l [y_3 + y_1 - 2y_2] + \text{terms } x \cdot y$$

$$\Rightarrow PE \propto \frac{1}{2} k_y (y_3 + y_1 - 2y_2)^2$$

Overall: $T = \frac{1}{2} [M \dot{x}_1^2 + m \dot{x}_2^2 + M \dot{x}_3^2 + M \dot{y}_1^2 + m \dot{y}_2^2 + M \dot{y}_3^2]$

$$U = \frac{1}{2} k_x [-(x_2 - x_1)^2 - (x_3 - x_2)^2] + \frac{1}{2} k_y [y_3 + y_1 - 2y_2]^2$$

$$U = \frac{1}{2} k_x \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} + \frac{1}{2} k_y \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} \begin{bmatrix} 1 & -2 & 1 \\ -2 & 4 & -2 \\ 1 & -2 & 1 \end{bmatrix} \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix}$$

$\frac{\partial U}{\partial x_i \partial x_j}$ $\frac{\partial U}{\partial y_i \partial y_j}$

Note: I've written x & y as separate vectors/matrices but the plan is to put together a big 6x6 matrix

$$U = \frac{1}{2} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ \dots \\ y_1 \\ y_2 \\ y_3 \end{pmatrix} \begin{pmatrix} k_x \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix} & 0 \\ \dots & \dots \\ 0 & k_y \begin{bmatrix} 1 & -2 & 1 \\ -2 & 4 & -2 \\ 1 & -2 & 1 \end{bmatrix} \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ \dots \\ y_1 \\ y_2 \\ y_3 \end{pmatrix}$$

But we need to do something else first.

```

1 U=kx/2( (q[2]-q[1])^2+(q[3]-q[2])^2)+ky/2(q[6]+q[4]-2 q[5])^2
2 U2=U /. {q[2]->Sqrt[M/m] q[2],q[5]->Sqrt[M/m] q[5]}
3 k=Table[Table[ D[U2,q[i],q[j]],{j,1,6}],{i,1,6}]
4 eigen=Eigensystem[k]

```

Note to get actual ω^2 divide by M

```

5
6 Out[4]= {{0, 0, 0, kx, -----, -----}, cut

```

undo scale change usual new bending mode

```

9 tx=IdentityMatrix[6]
10 tx[[2,2]]=Sqrt[M/m]
11 tx[[5,5]]=Sqrt[M/m]

```

```

12
13 Last[eigen].tx
14
15 Out[8]= {{0, 0, 0, -1, 0, 1}, {0, 0, 0, 2 Sqrt[-], Sqrt[-], 0},

```

```

16
17
18
19 > {1, 1, 1, 0, 0, 0}, {-1, 0, 1, 0, 0, 0}, {1, -----, 1, 0, 0, 0},

```

```

20
21
22
23 > {0, 0, 0, 1, -----, 1}

```

```

24
25
26 CO2
27 2438.1 cm^-1
28 1373.01
29 641.49
30 * ? 1+32/12
31 3.666666666666667
32 * ? (2438/1373)^2
33 3.153017114478638
34
35
36 Ge O2
37 1061.6
38 870.1
39 195.5
40
41 ? 1+32/72.6
42 1.440771349862259
43 * ? (1061.6/870.1)^2
44 1.488618741515112
45

```

The only pleasant way to solve 6x6 matrix problems is with Mathematica (but see below). Mathematica is willing & able to solve Eigensystem problems - but this is not exactly an eigensystem because the mass matrix $\begin{bmatrix} M & & & & & \\ & m & & & & \\ & & M & & & \\ & & & m & & \\ & & & & M & \\ & & & & & m \end{bmatrix}$ is not a scalar multiple of the unit matrix.

Solution - make a change in variables to make the mass matrix = $M \mathbb{1}$ Let $x_2 = \sqrt{\frac{M}{m}} q_2$ so $\frac{1}{2} m \dot{x}_2^2 = \frac{1}{2} M \dot{q}_2^2$

$\dot{y}_2 = \sqrt{\frac{M}{m}} \dot{q}_5 \Rightarrow (x_i \& y_i) \rightarrow q_i$

Make above substitutions into PE then take $\frac{\partial^2 L}{\partial q_i \partial q_j}$ to find the K matrix ... $L = \frac{1}{2} M \dot{\vec{q}} \mathbb{1} \dot{\vec{q}} - \frac{1}{2} \vec{q} [K] \vec{q}$

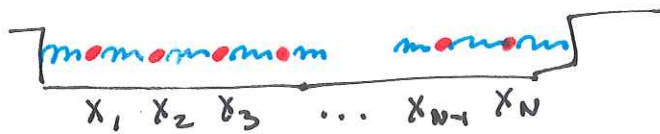
Lagrange $\Rightarrow M \ddot{\vec{q}} = - [K] \vec{q}$ try $\vec{q} = \vec{a} e^{i\omega t}$ constant

$$M \omega^2 \vec{a} = [K] \vec{a}$$

→ exactly an eigenvalue of [K]

Note: to get the physical locations $x \& y$: $x_2 = \sqrt{\frac{M}{m}} q_2$ etc

Now let's solve an arbitrarily large matrix problem by hand.



N masses m connected with identical springs

$$T = \frac{1}{2} \sum_{j=1}^N m \dot{x}_j^2$$

$$U = \frac{1}{2} k \left[\sum_{j=0}^N (x_{j+1} - x_j)^2 \right]$$

mass matrix = $m \mathbf{1}$

where $x_0 = 0$

$x_{N+1} = 0$ [from offset]

Note every x_j (except 0 & $N+1$) appears twice

$$(x_{j+1} - x_j)^2 + (x_j - x_{j-1})^2$$

mixed partial $\frac{\partial}{\partial x_j \partial x_{j+1}} \leftarrow \frac{\partial}{\partial x_j} \Rightarrow 2(x_{j+1} - x_j)(-1) + 2(x_j - x_{j-1})$

with either $j+1$ or $j-1$
 $\Rightarrow -2$

$$\frac{\partial^2}{\partial x_j^2} \Rightarrow 2+2=4$$

just off diagonal = -1

Note overall factor of $\frac{1}{2} k \Rightarrow$

$$[K] = k \begin{bmatrix} -1 & 2 & -1 & & 0 \\ & -1 & 2 & -1 & \\ 0 & & -1 & 2 & -1 \\ & & & -1 & 2 & -1 \\ & & & & & -1 & 2 & -1 \\ & & & & & & & -1 & 2 & -1 \end{bmatrix}$$

main diagonal has 2

Lagrange $\Rightarrow m \mathbf{1} \ddot{x} = -k [K] x$

Try $x = \vec{a} e^{i\omega t}$ $\Rightarrow m\omega^2 \vec{a} = k [K] \vec{a}$

since $\frac{k}{m}$ is usual $\omega_0^2 \Rightarrow \frac{\omega^2}{\omega_0^2} = \lambda = \text{eigenvalue}$

$$\lambda \vec{a} = [K] \vec{a}$$

$$\lambda a_j = -a_{j-1} + 2a_j - a_{j+1} \leftarrow \text{linear}$$

Try $a_j = \text{Re} [A e^{ij\gamma}]$ $\leftarrow \text{new factor}$

$$\Rightarrow \lambda = -e^{-i\gamma} + 2 - e^{i\gamma} \leftarrow A e^{ij\gamma} \text{ is common to all}$$

$$= 2(1 - \cos \gamma) = 4 \sin^2\left(\frac{\gamma}{2}\right)$$

$$L = 2\epsilon = 2\frac{\gamma}{2}$$

BC are going to tell us what γ are possible!

$$x_0 = \text{Re}[A e^{i0}] = 0 \Rightarrow A = \text{pure imaginary, eg } i$$

$$x_{N+1} = \text{Re}[A e^{i(N+1)\gamma}] = 0 \Rightarrow e^{i(N+1)\gamma} \text{ pure real}$$

only possible if $(N+1)\gamma = s\pi$
integer \rightarrow

$$\therefore \gamma = \frac{s\pi}{N+1} \quad s = 1, 2, \dots, N$$

if $s=1$ $j\gamma = \frac{j}{N+1}\pi$ will go $0 \rightarrow \pi$ as j goes $1 \rightarrow N$

$\frac{1}{2}$ cycle: 

if $s=2$ $j\gamma = \frac{j}{N+1}2\pi$ will go $0 \rightarrow 2\pi$ as j goes $1 \rightarrow N$


full cycle 

if $s=3$ $j\gamma = \frac{j}{N+1}3\pi$ will go $0 \rightarrow 3\pi$ as j goes $1 \rightarrow N$



conclude s is # half cycles as j goes $1 \rightarrow N$

if $s=N$ $\gamma = \frac{N}{N+1}\pi \approx \pi$ so $e^{ij\gamma}$ will alternate.


etc

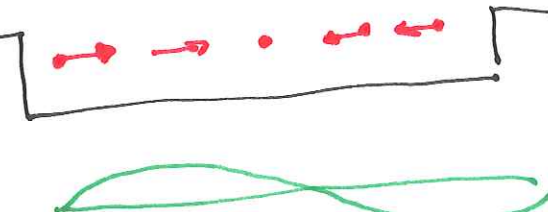
$s=1$: 

$$\lambda = \frac{v^2}{\omega^2} = 4 \sin^2\left(\frac{\gamma}{2}\right)$$

$$\approx \gamma^2 = \left[\frac{s\pi}{N+1}\right]^2$$

$$\omega = \omega_0 \left[\frac{s\pi}{N+1}\right]$$

much lower freq than ω_0 cuz adjacent particles move together

$s=2$: 

I hope this reminds you of resonance in a closed-closed organ pipe because that's exactly what it is.