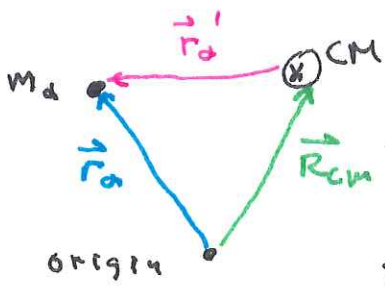


Chapter 10 starts with a review of

Composite this.

Useful Lemma:

$$\sum m_d \vec{r}'_d = 0$$



$$\vec{R}_{CM} = \frac{1}{M} \sum m_d \vec{r}_d$$

$$\vec{V}_{CM} = \frac{1}{M} \sum m_d \vec{v}_d$$

$$\vec{A}_{CM} = \frac{1}{M} \sum m_d \vec{a}_d$$

Because of Newton 3rd, internal forces cancel $\dot{\vec{F}}_{ext} = M \vec{A}$

ie the source of the force is a particle in the system

$$KE: T = \sum \frac{1}{2} m_d v_d^2 = \underbrace{\frac{1}{2} M V_{CM}^2}_{\text{"OF" CM}} + \underbrace{\sum \frac{1}{2} m_d v_d'^2}_{\text{"ABOUT" CM}}$$

FYJ: the test book has not stressed that for uniform gravity the net force can be thought of as applied at CM

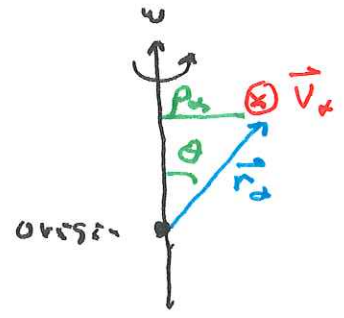
$$PE - PE = \sum m_d g z_d = M g z_{CM}$$

$$\text{torque: } \sum \vec{r}_d \times m_d \vec{g} = M \vec{R}_{CM} \times \vec{g}$$

For rotating rigid bodies: $\vec{v}_d = \vec{\omega} \times \vec{r}_d$

$$|\vec{v}_d| = |\omega| |\vec{r}_d| \sin \theta = \omega \rho_d$$

$$\text{So } T = \sum \frac{1}{2} m_d (\omega \rho_d)^2 = \frac{1}{2} \sum m_d \rho_d^2 \omega^2$$



Angular Momentum: $\vec{L} = \sum m_d \vec{r}_d \times \vec{v}_d$

$$= M \vec{R}_{CM} \times \vec{V}_{CM} + \sum m_d \vec{r}'_d \times \vec{v}'_d$$

"OF" CM
Orbital depends on origin

"ABOUT" CM
spin does not depend on origin

we called this moment of inertia I but its about to become more complex

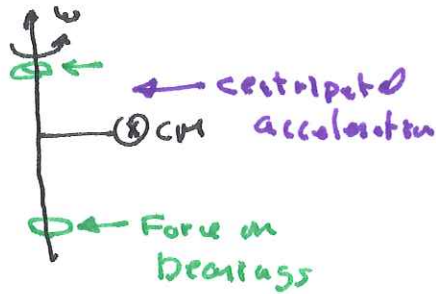
Note by above comments gravity ($m\vec{g}$) never has a torque about the CM

$$\frac{dL_{orbit}}{dt} = \underbrace{\vec{R}_{CM} \times \vec{F}_{ext}}_{\text{Torque OF CM}}$$

$$\frac{dL_{spin}}{dt} = \sum \underbrace{r'_d \times \vec{F}_{ext}}_{\text{Torque ABOUT CM}}$$

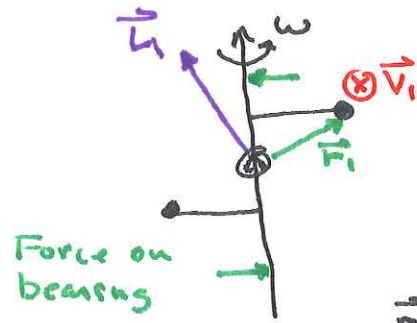
Balanced & Not rotation

Case of CM off axis:



So in general \vec{L} & $\vec{\omega}$ are not pointed in same direction so $\vec{L} = I \vec{\omega}$ not generally possible.

CM on axis but masses "unbalanced"



For other mass: $\vec{r}_2 = -\vec{r}_1$
 $\vec{v}_2 = -\vec{v}_1$
 so $\vec{L}_2 = \vec{L}_1$
 so \vec{L}_{total} must rotate
 $\Rightarrow \frac{d\vec{L}}{dt} \neq 0 \Rightarrow$ Torque required
 \vec{L} out-of-page

$$\vec{L} = \sum m_a \vec{r}_a \times \vec{v}_a = \sum m_a \vec{r}_a (\vec{\omega} \times \vec{r}_a)$$

$$= \sum m_a \left\{ \vec{\omega} r_a^2 - \underbrace{\vec{r}_a (\vec{\omega} \cdot \vec{r}_a)} \right\}$$

This term $\Rightarrow \sum m_a r_a^2 \vec{\omega}$
 so result is in $\vec{\omega}$ direction

This is "problem" term.
 More generally: $\vec{A}(\vec{B} \cdot \vec{C})$

weights as matrix:

$$\begin{pmatrix} A_x(B_x C_x + B_y C_y + B_z C_z) \\ A_y(\dots) \\ A_z(\dots) \end{pmatrix}$$

$$\begin{bmatrix} A_x B_x & A_x B_y & A_x B_z \\ A_y B_x & A_y B_y & A_y B_z \\ A_z B_x & A_z B_y & A_z B_z \end{bmatrix} \begin{bmatrix} C_x \\ C_y \\ C_z \end{bmatrix} = (\vec{A} \vec{B}) \cdot \vec{C}$$

"Dyadic" $\vec{A} \vec{B}$

No dot or cross here - just juxtaposed

here $\vec{A} = \vec{r}_a$ & $\vec{B} = \vec{r}_a$

so

$$\begin{bmatrix} xx & xy & xz \\ yx & yy & yz \\ zx & zy & zz \end{bmatrix}$$

$$L = \left\{ \begin{bmatrix} \sum m_d r_d^2 & 0 & 0 \\ 0 & \sum m_d r_d^2 & 0 \\ 0 & 0 & \sum m_d r_d^2 \end{bmatrix} - \begin{bmatrix} \sum m_d x_d y_d & \sum m_d x_d z_d & \sum m_d y_d z_d \\ \sum m_d y_d x_d & \sum m_d y_d z_d & \sum m_d y_d z_d \\ \sum m_d z_d x_d & \sum m_d z_d y_d & \sum m_d z_d z_d \end{bmatrix} \right\}$$

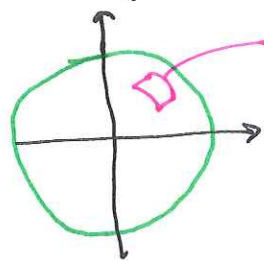
$$= \begin{bmatrix} \sum m_d (r_d^2 - x_d^2) & -\sum m_d x_d y_d & -\sum m_d x_d z_d \\ -\sum m_d y_d x_d & \sum m_d (r_d^2 - y_d^2) & -\sum m_d y_d z_d \\ -\sum m_d z_d x_d & -\sum m_d z_d y_d & \sum m_d (r_d^2 - z_d^2) \end{bmatrix} \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix}$$

$$= \vec{I} \cdot \vec{\omega}$$

↑
moment of inertia tensor

↑
 $= x_d^2 + y_d^2$
↑
re everything except z

Examples — Disk — "lamina" — thin sheet — all mass at $z=0$ plane



$$dA = dr r d\phi$$

$$x = r \cos \phi$$

$$y = r \sin \phi$$

$$z = 0$$

$$\rightarrow dm = \sigma dA$$

$$\frac{\text{mass}}{\text{Area}} = \sigma = \frac{M}{\pi R^2}$$

$$I_{xy} = - \int \sigma dr r d\phi xy = -\sigma \int r^2 \sin \phi \cos \phi r dr d\phi$$

$$I_{xx} = \int \sigma dr r d\phi (y^2 + z^2) = \sigma \int r^2 \sin^2 \phi r dr d\phi$$

$$= \frac{M}{\pi R^2} \frac{1}{4} R^4 \pi = \frac{1}{4} MR^2$$

Note: $I_{yy} \rightarrow \int \cos^2 \phi d\phi = \frac{1}{2} \cdot 2\pi$ so same result

$$I_{zz} = \int \sigma dr r d\phi (x^2 + y^2) = I_{xx} + I_{yy} = \frac{1}{2} MR^2$$

$$\vec{I} = \frac{1}{2} MR^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

↑
see that this is a general result for lamina

Cube at CM:



take side as $2s$
integrals: \int_{-s}^s

$$I_{xy} = - \int \rho \, dx \, dy \, dz \, (xy)$$

$$\int_{-s}^s y \, dy = \left. \frac{1}{2} y^2 \right|_{-s}^s = 0$$

$$I_{xx} = \int \rho \, dx \, dy \, dz \, (y^2 + z^2)$$

z will work just like y

$$\frac{M}{(2s)^3}$$

$$\int_{-s}^s dx \int_{-s}^s dz \int_{-s}^s y^2 dy$$

\downarrow $2s$ \downarrow $2s$ \downarrow $\frac{2s^3}{3}$

$$= \rho \cdot 2s \cdot 2s \cdot \frac{2}{3} s^3 \times 2$$

$$= \frac{2}{3} M s^2$$

$$\hat{I} = \frac{2}{3} M s^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

so here \vec{L} & $\vec{\omega}$ point same direction