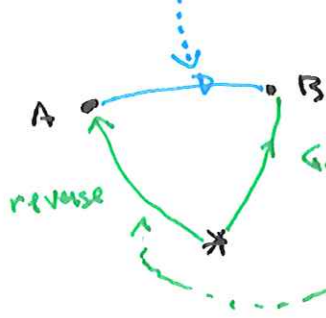


Now define βE $U(\vec{r}) = \int_{*}^{\vec{r}} \vec{F} \cdot d\vec{r}$ where \vec{F} is assumed conservative - i.e. this integral does not depend on path

Note: $\int_A^B \vec{F} \cdot d\vec{r} = T(B) - T(A)$



arbitrary starting point. If changed just adds a constant to U which has no effect.

$$-\int_{*}^A \vec{F} \cdot d\vec{r} + \int_{*}^B \vec{F} \cdot d\vec{r} = -U(B) + U(A)$$

SO: $T(B) - T(A) = -U(B) + U(A)$

$$T(B) + U(B) = T(A) + U(A)$$

Conservation of Energy

Remark: If one of the forces is non conservative single that one out from

Integral: $\int_A^B \vec{F} \cdot d\vec{r} = \underbrace{-U(B) + U(A)}_{\text{work done by conservative forces}} + \underbrace{\int_A^B \vec{F}_{nc} \cdot d\vec{r}}_{\text{work done by non conservative forces}}$

$T(B) - T(A)$

$$\Rightarrow (T(B) + U(B)) - (T(A) + U(A)) = \int_A^B \vec{F}_{nc} \cdot d\vec{r}$$

Work Energy Thm

: work done by non conservative forces results in changed energy

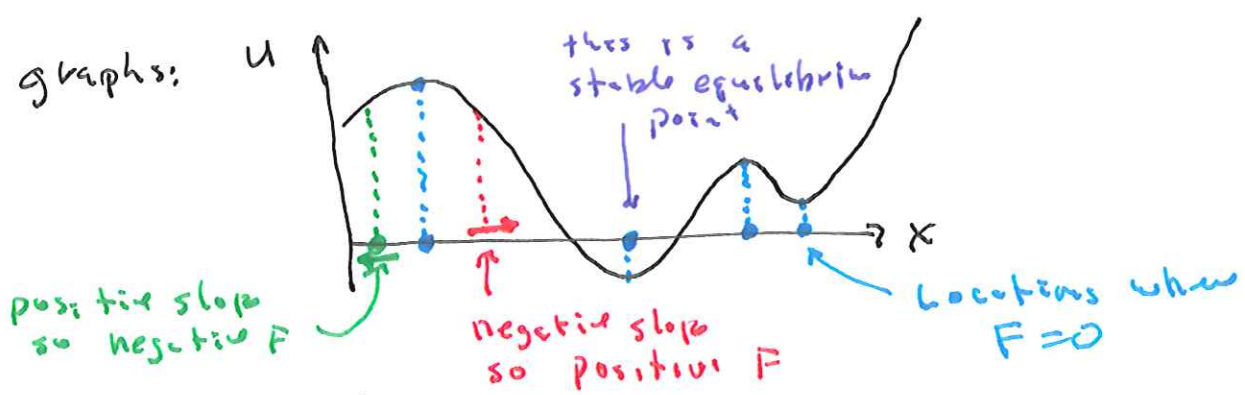
1-d case: If $F(x)$ just depends on x its automatically conservative as there really aren't different paths

in 1-d: $U(x) = -\int_{*}^x F(x) dx$ $-\frac{dU}{dx} = F_x$

Eg SHO: $F = -kx \Rightarrow U(x) = \int_{*}^x kx dx = \frac{1}{2} kx^2$

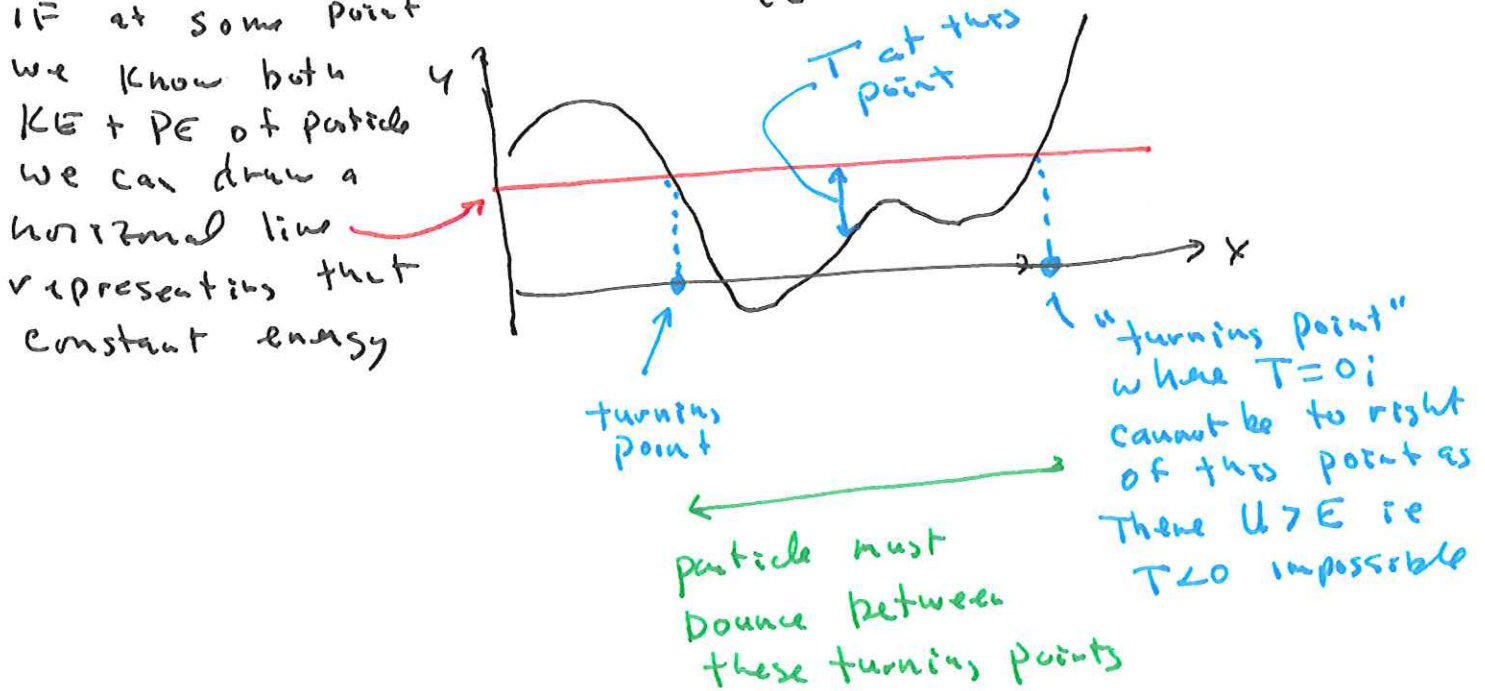
unimportant constant usually taken to be zero

PE graphs:



see is unstable - if right on blue dot $F=0$ so could remain at rest but is moved a bit to right red force pushes you more to right. If moved a bit to left green force moves you further to left

If at some point we know both KE + PE of particle we can draw a horizontal line representing that constant energy



Instead of using the 2nd order diff eg: $F(x) = m\ddot{x}$

we can use the equivalent 1st order diff eg - Conservation of E

$$\frac{1}{2} m \dot{x}^2 + U(x) = E \Rightarrow \dot{x} = \sqrt{\frac{2}{m} (E - U(x))}$$

$$F(x_f) - F(x_i) = \cancel{v_f v_i} = \int_{x_i}^{x_f} \frac{dx}{\sqrt{\frac{2}{m} (E - U(x))}} = \int dt = t_f - t_i$$

↑
anti derivative
from that integral

Eg SHO: $\frac{1}{2} m \dot{x}^2 + \frac{1}{2} k x^2 = E$ ← we write this constant in a more convenient way
 $= \frac{1}{2} k A^2$ ← new constant

$$\dot{x} = \sqrt{\frac{k}{m} (A^2 - x^2)}$$

Call this new constant ω

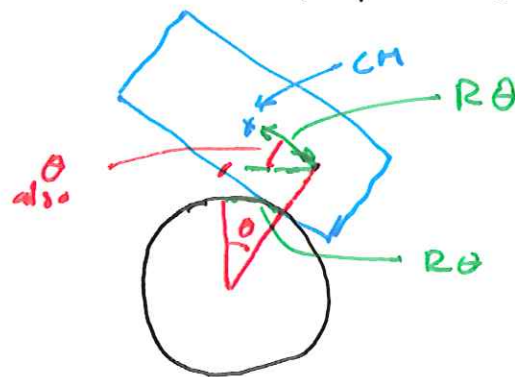
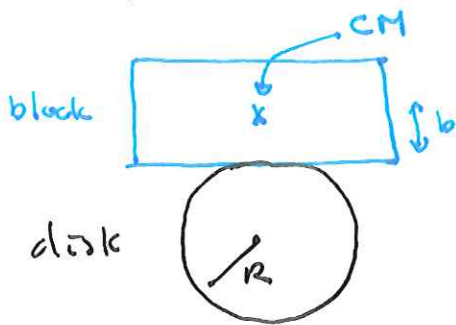
$$\frac{dx}{\sqrt{A^2 - x^2}} = \omega dt$$

$$\int \frac{dx}{\sqrt{A^2 - x^2}} = \sin^{-1}\left(\frac{x}{A}\right) = \omega(t - t_0)$$

$$\frac{x}{A} = \sin(\omega(t - t_0))$$

$$x = A \sin(\omega(t - t_0))$$

Eg → Finding U in a more complex case — stable or unstable



Red vector = $((R+b) \sin \theta, (R+b) \cos \theta)$

Green vector = $(-R\theta \cos \theta, R\theta \sin \theta)$

CM location = $((R+b) \sin \theta - R\theta \cos \theta, (R+b) \cos \theta + R\theta \sin \theta)$

Gravitational PE = mgy

$$U = mg \left[\underbrace{(R+b) \cos \theta}_{\cos \theta \approx 1 - \frac{1}{2!} \theta^2 + \frac{1}{4!} \theta^4 - \dots} + \underbrace{R\theta \sin \theta}_{\sin \theta \approx \theta - \frac{1}{3!} \theta^3 + \frac{1}{5!} \theta^5 - \dots} \right]$$

$$= mg \left[(R+b) - \underbrace{\left(\frac{R+b}{2} \right) \theta^2}_{\left(\frac{R-b}{2} \right) \theta^2 \text{ if } b > R \text{ unstable}} + R\theta + \text{terms } \theta^4, \theta^6 \dots \right]$$