

Circular Orbits (as in |a|)

$$F = -\frac{\alpha}{r^2} = -\mu \omega^2 r$$

centrifugal acceleration

both force and acceleration point toward attracting mass

$$\alpha = G m_1 m_2$$

$$\mu = \frac{m_1 m_2}{m_1 + m_2}$$

$$\frac{\alpha}{\mu} = G(m_1 + m_2)$$

$$\frac{\alpha}{\mu r^3} = \omega^2$$

$$KE = \frac{1}{2} \mu (\omega r)^2 = \frac{1}{2} \mu \frac{\alpha}{\mu r}$$

$$PE = -\frac{\alpha}{r}$$

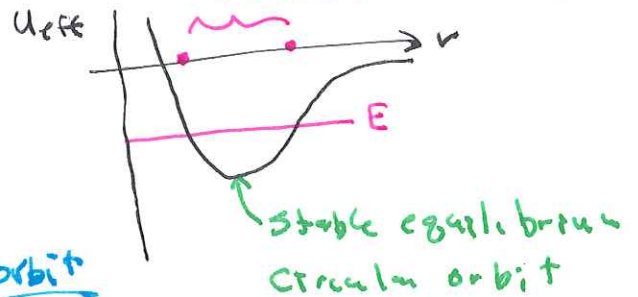
$$E = KE + PE = -\frac{\alpha}{2r}$$

$$L = \mu r^2 \omega = \sqrt{\mu \alpha r^3}$$

Lagrange: $U_{eff} = \frac{L^2}{2\mu r^2} + U$

centrifugal potential = transverse KE

For small r this $\rightarrow \infty$
 For larger this $\rightarrow 0^-$
 turning pts: r_{min}/r_{max}



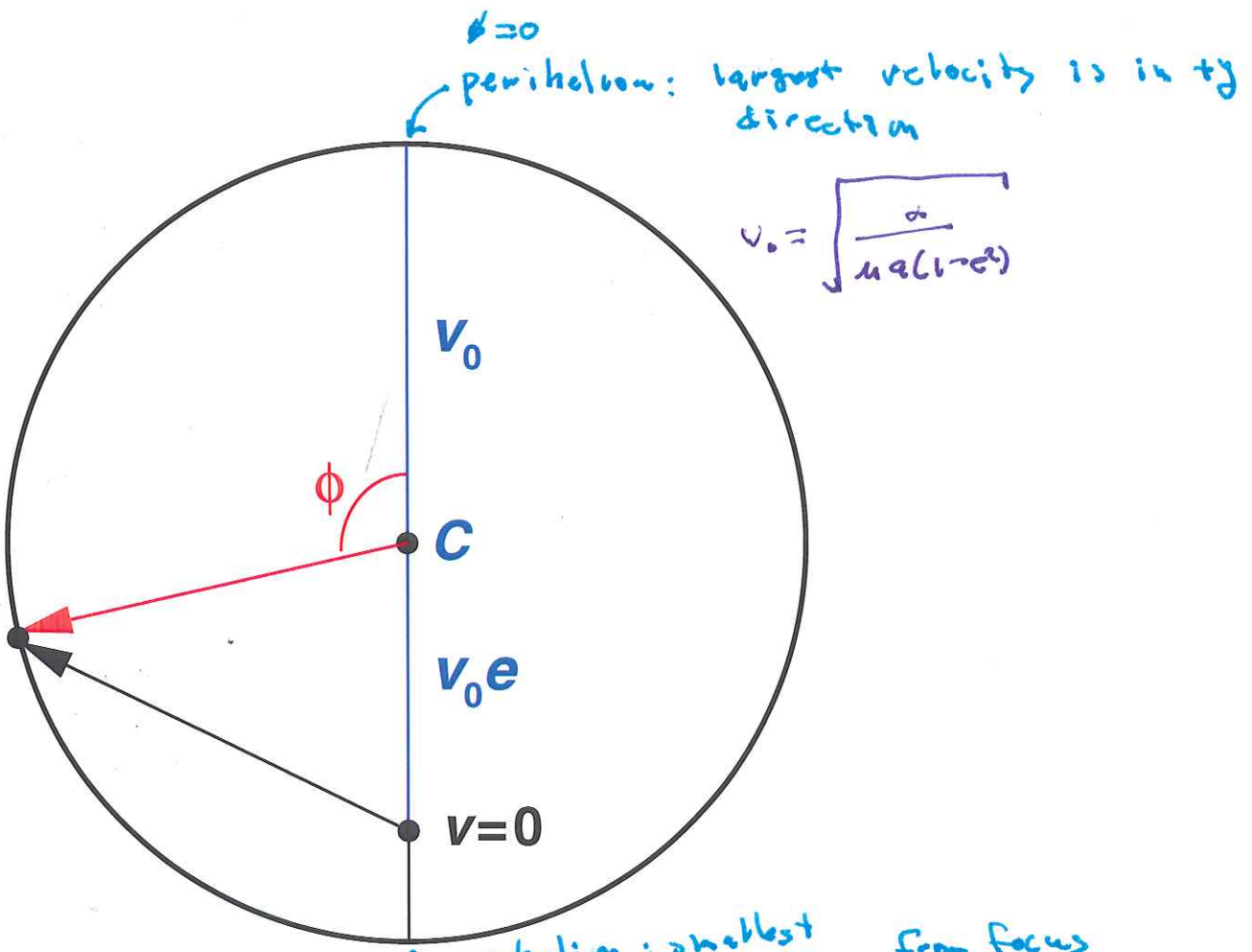
\rightarrow show in textbook: seek orbit

ie $r(\phi)$, switch variables $u = \frac{1}{r}$ seek $u(\phi)$

produces a solvable diff eq $\rightarrow r = \frac{a(1-e^2)}{1+e \cos \phi}$ ← ellipse

using graphical approach: $\frac{d\vec{v}}{dt}$ points in direction $\phi + 180^\circ$
 ie in $-\vec{r}$ direction (good - that's direction of force)

$$|\frac{d\vec{v}}{dt}| = v_0 \dot{\phi} = v_0 \frac{L}{\mu r^2} = \sqrt{\frac{\alpha}{\mu a(1-e^2)}} \frac{|\alpha \mu a(1-e^2)|}{\mu r^2} = \frac{\alpha}{\mu r^2} \checkmark$$



from center

eccentric anomaly

$$\vec{r} = (a \cos u, b \sin u)$$

$$b = a \sqrt{1-e^2}$$

$$\tan \frac{\phi}{2} = \sqrt{\frac{1+e}{1-e}} \tan \frac{u}{2}$$

$$a - e \sin u = \frac{at}{v}$$

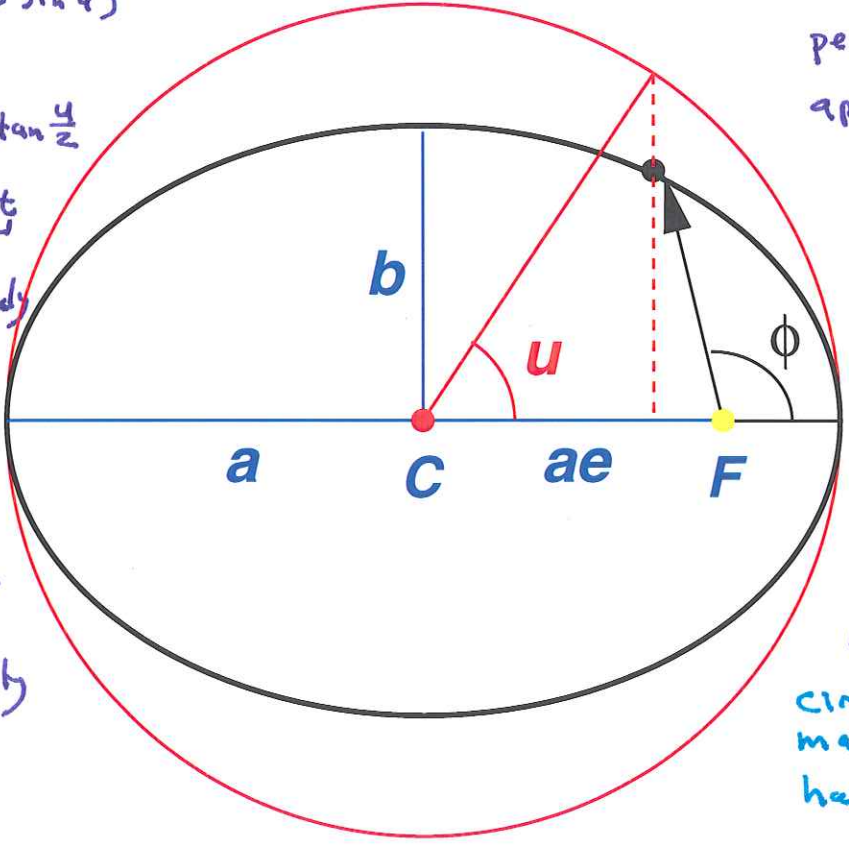
mean anomaly

ϕ = true anomaly

a = semi major axis

b = semi minor axis

e = eccentricity



from focus

$$r = \frac{a(1-e^2)}{1+e \cos \phi}$$

$$\text{peri} = r_{\min} = a(1-e)$$

$$\text{apo} = r_{\max} = a(1+e)$$

$$w = \sqrt{\frac{\alpha}{\mu r^3}}$$

$$E = \frac{-\alpha}{2a}$$

$$L = \sqrt{\alpha \mu a(1-e^2)}$$

do not depend on eccentricity

circular orbits ($e=0$) have max L; line orbits ($e=1$) have min L