

Digital Filters

Finite Impulse Response (non recursive)

response is for a finite time

→ input is zero except at a single time

$$\text{output } y_k = a_0 x_k + a_1 x_{k-1} + a_2 x_{k-2} + \dots + a_n x_{k-n}$$

eg "running average" where a_i are constant = $\frac{1}{n+1}$

Note: seek response to a unit step & a unit sinusoidal at some freq ω

Infinite Impulse Response (recursive)

$$\text{output } y_k = a_0 x_k + a_1 x_{k-1} + \dots + a_n x_{k-n} \\ - (b_1 y_{k-1} + b_2 y_{k-2} + \dots + b_m y_{k-m})$$

Just as a damped driven harmonic oscillator has a homogeneous solution (that decays & depends on initial conditions) and a steady particular solution that does not depend on initial conditions same applies here to digital filters - if $x_k = e^{i\omega k}$ long term $y_k = A e^{i\omega k}$ where A will depend on the frequency.

$$A (e^{ik\omega} + b_1 e^{i(k-1)\omega} + b_2 e^{i(k-2)\omega} + \dots + b_m e^{i(k-m)\omega})$$

$$= a_0 e^{ik\omega} + a_1 e^{i(k-1)\omega} + a_2 e^{i(k-2)\omega} + \dots + a_n e^{i(k-n)\omega}$$

$$\Rightarrow A = \frac{a_0 + a_1 e^{-i\omega} + \dots + a_n e^{-in\omega}}{1 + b_1 e^{-i\omega} + b_2 e^{-2i\omega} + \dots + b_m e^{-mi\omega}}$$

so can use above to find amplitude & phase of output.

Z-transform: a way to deal with entire sequence $\{h_k\}$ at once by making it into a function

$$H(z) \equiv \sum_{k=0}^{\infty} \frac{h_k}{z^k}$$

Remark: this is formal process rather than an analytical process so we're not really concerned about things like convergence, but note that the ratio test \Rightarrow

$$\frac{|h_{k+1}/h_k|}{|z|} \quad \text{so if } \left| \frac{h_{k+1}}{h_k} \right| \text{ is bounded we'll}$$

have convergence for sufficiently large $|z|$. Typically

$$\left| \frac{h_{k+1}}{h_k} \right| \leq 1 \quad \text{so we'll have convergence for } |z| > 1.$$

This can be a bit tricky as we will often set $z = e^{i\omega}$ where then $|z| = 1$.

Examples:

impulse: $h_k = \{1, 0, 0, 0, \dots\}$

$$H(z) = 1$$

step: $h_k = \{1, 1, 1, 1, \dots\}$

$$H(z) = \frac{1}{1 - 1/z}$$

decay: $h_k = \{e^{-\alpha k}\}$

$$H(z) = \frac{1}{1 - \frac{1}{e^\alpha} z}$$

oscillation: $h_k = \{e^{i\omega k}\}$

$$H(z) = \sum \left(\frac{e^{i\omega}}{z} \right)^k = \frac{1}{1 - \frac{e^{i\omega}}{z}} = \frac{1}{1 - \frac{e^{-i\omega}}{z}} \frac{1 - e^{-i\omega}/z}{1 - e^{i\omega}/z}$$

$$= \frac{1 - e^{-i\omega}/z}{1 - \frac{z}{2} \cos \omega + \frac{1}{2} z}$$

Note: if want Re or Im part of above, just do it.

Note: we move a h_k one step to future by multiplying by z^{-1} , thus a 3-step running average of h_k :

$$g_k = \frac{1}{3} (h_k + h_{k-1} + h_{k-2})$$

$$G(z) = \frac{1}{3} (1 + z^{-1} + z^{-2}) H(z)$$

Examples: difference: $g_k = h_k - h_{k-1}$

$$G(z) = (1 - z^{-1}) H(z)$$

inverse process:

$$g_k = \sum_{i=0}^k h_i$$

$$G(z) = \frac{1}{(1 - z^{-1})} H(z)$$

Remark: the above are inverses of each other — like integration & differentiation. Encryption/decryption seeks such pairs in digital encryption (where we assume lossless-perfect-transmission) huge scrambles are possible. In Analog encryption, decryption must not be sensitive to noise.

Remark: Literature nobel prize winner Aleksander Solzhenitsyn wrote about folks working in analog "scrambling" in Stalins Gulags in the novel "In the First Circle" refers to Dante's *Inferno*

Typical IIR filter:

$$g_k = \sum_{i=0}^M a_i h_{k-i} - \sum_{i=1}^M b_i g_{k-i}$$

$$G(z) = \frac{\sum a_i z^{-i}}{1 + \sum b_i z^{-i}} H(z)$$

transfer function

$$A(z)$$

Remark: the transfer function is the response to an impulse: $H(z)$
 Remark: since everything is linear we can determine result of an arbitrary input as a sequence of impulses

Remark: previous work - if $x_k = e^{i\omega k}$ output (eventually)

becomes $y_k = \underbrace{A(e^{i\omega})}_{\text{the magnitude at this } \omega \text{ plotted as our Bode plot.}} e^{i\omega k}$

Example: 2 order low pass Butterworth 2000/10000

$$A(z) = \frac{1 + 2z^{-1} + z^{-2}}{1 + 0.385z^{-1} + 0.1458z^{-2}}$$

the numerator factors: $(1 + z^{-1})^2$ so it has double zero @ $z = -1$ i.e. $z = e^{\pm i\pi}$ i.e. at Nyquist freq.

the denominator can be written $(1 - \frac{z_1}{z})(1 - \frac{z_1^*}{z})$
 where $z_1 = 0.18476 + 0.40209i = 0.4425 e^{i(1.14)}$
 "poles"
 $f = 1814 \text{ Hz} \leftarrow 2\pi \frac{f}{f_0}$

Remark: the IIR of a term $\frac{1}{(1 - \frac{z_1}{z})}$ is $z_1^k = r^k e^{i2\pi \frac{f}{f_0} k}$ i.e. decaying amplitude at $f = 1814 \text{ Hz}$.

with a little partial fractions work

$$A(z) = 5.1068 + \frac{-4.1068 + 3.887z^{-1}}{(1 - z_1/z)(1 - z_1^*/z)}$$

$$A + B = -4.1068$$

$$-(Az_1^k + Bz_1^k) = 3.887$$

$$= \frac{A}{(1 - z_1/z)} + \frac{B}{(1 - z_1^*/z)}$$

$$= \frac{-2.053 - 3.890i}{(1 - z_1/z)} + \frac{-2.053 + 3.890i}{(1 - z_1^*/z)}$$

Example: 2 order Butterworth high pass

3000/10000

$$A(z) = \frac{1 - 2z^{-1} + z^{-2}}{1 + .3695z^{-1} + .1958z^{-2}}$$

numerator: $(1 - z^{-1})^2$ so has double zero at $z=1$ i.e. e^{i0}

The denominator can be written $(1 - \frac{z_1}{z})(1 - \frac{z_1^*}{z})$

where $z_1 = -.18476 + .40209i = .4425 e^{i2.0015}$

$f = 3186 \leftarrow 2\pi \frac{f}{f_0}$

Partial Fraction expansion.

$$A(z) = 5.1068 - \frac{4.1068 + 3.587z^{-1}}{(1 - \frac{z_1}{z})(1 - \frac{z_1^*}{z})}$$

$$A + B = 4.1068$$

$$Az_1 + Bz_1 = -3.587$$

$$\frac{A}{(1 - z_1/z)} + \frac{B}{(1 - z_1^*/z)}$$

$$A(z) = 5.1068 + \frac{-2.053 + 32.59i}{(1 - z_1/z)} + \frac{-2.053 - 32.59i}{(1 - z_1^*/z)}$$

digital filter notation

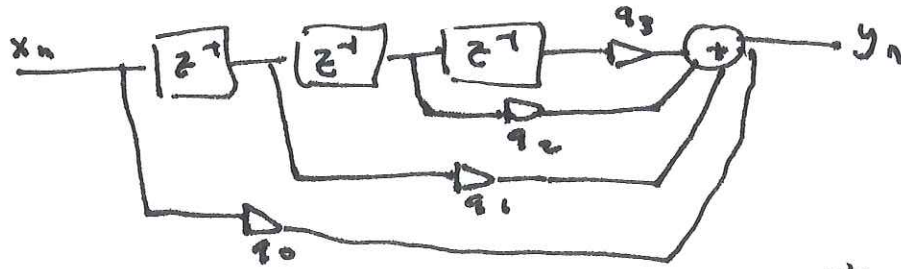
(+) adder

▷ multiplication by constant

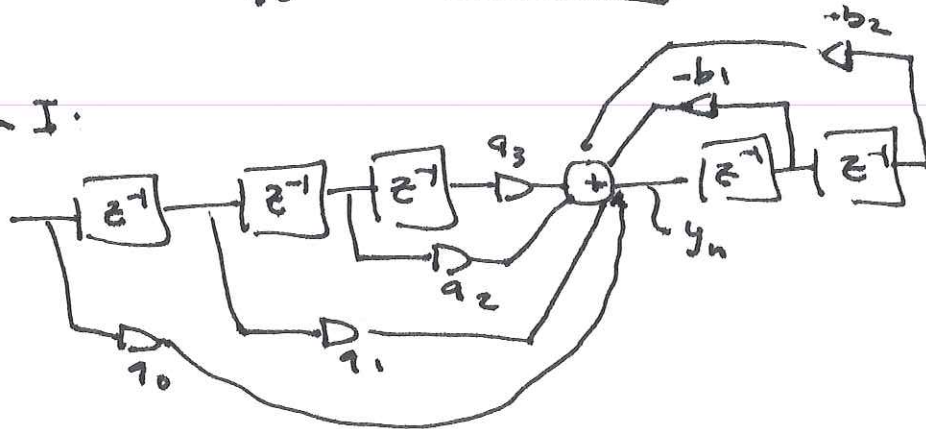
z^{-1}

delay 1 or previous

FIR:



IIR version I:



version II

Let $A(z) = \frac{N(z)}{D(z)}$

is consider $W(z) = \frac{1}{D(z)} X(z)$

the $Y(z) = N(z)W(z)$

