17-13. From graph: $\Delta p_m = 8 \times 10^{-3}$ Pa; $T = 2 \times 10^{-3}$ s. Data from problem: $\rho = 1.21 \text{ kg/m}^3$ (later $1.35 \,\mathrm{kg/m^3}$, $v = 343 \,\mathrm{m/s}$ (later $320 \,\mathrm{m/s}$)

(a) Eq. 17-15 reports: $s_m = \Delta p_m/v \rho \omega$ and as usual $\omega = 2\pi/T$ so:

$$
s_m = \frac{\Delta p_m}{v \rho \omega} = \frac{\Delta p_m T}{v \rho 2\pi} = \frac{8 \times 10^{-3} \cdot 2 \times 10^{-3}}{343 \cdot 1.21 \cdot 2\pi} \cdot \frac{N/m^2 \cdot s}{m/s \cdot kg/m^3} = 6.14 \times 10^{-9} \text{ m}
$$

(b)

$$
k = \frac{\omega}{v} = \frac{2\pi}{Tv} = \frac{2\pi}{2 \times 10^{-3} \text{ s} \cdot 343 \text{ m/s}} = 9.16 \text{ m}^{-1}
$$

(c)

$$
\omega = \frac{2\pi}{T} = \frac{2\pi}{2 \times 10^{-3} \,\mathrm{s}} = 3140 \,\mathrm{s}^{-1}
$$

(d)

$$
s_m = \frac{\Delta p_m T}{v \rho 2\pi} = \frac{8 \times 10^{-3} \cdot 2 \times 10^{-3}}{320 \cdot 1.35 \cdot 2\pi} \frac{N/m^2 \cdot s}{m/s \cdot kg/m^3} = 5.89 \times 10^{-9} \text{ m}
$$

(e)

$$
k = \frac{2\pi}{Tv} = \frac{2\pi}{2 \times 10^{-3} \text{ s} \cdot 320 \text{ m/s}} = 9.82 \text{ m}^{-1}
$$

(f) unchanged

- 17-17. Let L_1 be the distance from the closer speaker to the listener. The distance from the other speaker to the listener is $L_2 = \sqrt{L_1^2 + d^2}$, where d is the distance between the speakers. The phase difference at the listener is $\phi = 2\pi (L_2 - L_1)/\lambda$, where λ is the wavelength.
	- (a) For a minimum in intensity at the listener, $\phi = (2n+1)\pi$, where n is an integer. Thus $\lambda =$ $2(L_2 - L_1)/(2n + 1)$. The frequency is

$$
f = \frac{v}{\lambda} = \frac{(2n+1)v}{2(\sqrt{L_1^2 + d^2} - L_1)} = \frac{(2n+1)(343 \text{ m/s})}{2(\sqrt{(3.75 \text{ m})^2 + (2.00 \text{ m})^2} - 3.75 \text{ m})} = (2n+1)(343 \text{ Hz}).
$$

- (b) 3
- (c) 5
- (d) For a maximum in intensity at the listener, $\phi = 2n\pi$, where n is any positive integer. Thus $\lambda = (1/n) \left(\sqrt{L_1^2 + d^2} - L_1 \right)$ and

$$
f = \frac{v}{\lambda} = \frac{nv}{\sqrt{L_1^2 + d^2} - L_1} = \frac{n(343 \text{ m/s})}{\sqrt{(3.75 \text{ m})^2 + (2.00 \text{ m})^2} - 3.75 \text{ m}} = n(686 \text{ Hz}) .
$$

Since $20,000/686 = 29.2$, n must be in the range from 1 to 29 for the frequency to be audible and $f = 686, 1372, \ldots, 19890$ Hz.

- (e) 2
- (f) 3

17-36. At the beginning of the exercises and problems section in the textbook, we are told to assume $v_{\rm sound} =$ 343 m/s unless told otherwise. The second harmonic of pipe A is found from Eq. 17-39 with $n = 2$ and $L = L_A$, and the third harmonic of pipe B is found from Eq. 17-41 with $n = 3$ and $L = L_B$. Since these frequencies are equal, we have

$$
\frac{2v_{\text{sound}}}{2L_A} = \frac{3v_{\text{sound}}}{4L_B} \implies L_B = \frac{3}{4}L_A.
$$

- (a) Since the fundamental frequency for pipe A is 300 Hz, we immediately know that the second harmonic has $f = 2(300) = 600$ Hz. Using this, Eq. 17-39 gives $L_A = (2)(343)/2(600) = 0.572$ m.
- (b) The length of pipe B is $L_B = \frac{3}{4}L_A = 0.429$ m.
- 17-47. Each wire is vibrating in its fundamental mode so the wavelength is twice the length of the wire $(\lambda = 2L)$ and the frequency is $f = v/\lambda = (1/2L)\sqrt{\tau/\mu}$, where $v = \sqrt{\tau/\mu}$ is the wave speed for the wire, τ is the tension in the wire, and μ is the linear mass density of the wire. Suppose the tension in one wire is τ and the oscillation frequency of that wire is f_1 . The tension in the other wire is $\tau + \Delta \tau$ and its frequency is f_2 . You want to calculate $\Delta \tau / \tau$ for $f_1 = 600$ Hz and $f_2 = 606$ Hz. Now, $f_1 = (1/2L)\sqrt{\tau/\mu}$ and $f_2 = (1/2L)\sqrt{(\tau + \Delta \tau)/\mu}$, so

$$
f_2/f_1 = \sqrt{(\tau + \Delta \tau)/\tau} = \sqrt{1 + (\Delta \tau/\tau)}.
$$

This leads to

$$
\Delta \tau / \tau = (f_2/f_1)^2 - 1 = [(606 \text{ Hz})/(600 \text{ Hz})]^2 - 1 = 0.0201.
$$

17-69. Since the amplitudes are all the same: \sum 4 $k=1$ $a_k e^{i \delta_k} = a_1 \sum$ 4 $k=1$ $e^{i\delta_k}$, however the sum of the phase shifts is exactly zero: $e^{0i} + e^{.7\pi i} + e^{\pi i} + e^{1.7\pi i} = 1 + e^{\pi i} + e^{.7\pi i} (1 + e^{\pi i}) = 0$. As usual we have used:

$$
h(t) = a_1 \cos(\omega t + \delta_1) + a_2 \cos(\omega t + \delta_2) + \dots + a_N \cos(\omega t + \delta_N)
$$

\n
$$
= \sum_{k=1}^N a_k \cos(\omega t + \delta_k)
$$

\n
$$
= \text{Re}\left[(a_1 e^{i\delta_1} + a_2 e^{i\delta_1} + \dots + a_N e^{i\delta_N}) e^{i\omega t} \right]
$$

\n
$$
= \text{Re}\left[\left(\sum_{k=1}^N a_k e^{i\delta_k} \right) e^{i\omega t} \right]
$$

17-27. (a) From Eq. 17-29 we have:

$$
10 \log \left(\frac{I}{I_0}\right) = \beta
$$

$$
\log \left(\frac{I}{I_0}\right) = \beta/10
$$

$$
\left(\frac{I}{I_0}\right) = 10^{\beta/10}
$$

$$
I = I_0 10^{\beta/10}
$$

$$
= 10^{-12} 10^{70/10} = 10^{-5} W/m^2
$$

(b) Note that intensity ratios are related to dB differences:

$$
10 \log \left(\frac{I_1}{I_0}\right) = \beta_1
$$

$$
10 \log \left(\frac{I_2}{I_0}\right) = \beta_2
$$

since subtracting these two equations shows:

$$
10\left[\log\left(\frac{I_1}{I_0}\right) - \log\left(\frac{I_2}{I_0}\right)\right] = 10\left[\log\left(\frac{I_1}{I_2}\right)\right] = \beta_1 - \beta_2
$$

so

$$
\left(\frac{I_1}{I_2}\right) = 10^{(\beta_1 - \beta_2)/10}
$$

Here final/initial is requested:

$$
\left(\frac{I_f}{I_i}\right) = 10^{(\beta_f - \beta_i)/10} = 10^{(50 - 70)/10} = 10^{-2}
$$

(c) Using Eq. 17-27 we can realted intensities to displacement amplitudes:

$$
s_m = \sqrt{\frac{2I}{\rho v \omega^2}} = \sqrt{\frac{2 \cdot 10^{-5}}{1.21 \cdot 343 \cdot (2\pi \, 500)^2}} = 69.9 \times 10^{-9} \, \text{m}
$$

Unit check:

$$
\sqrt{\frac{(N\cdot m/s)\cdot m^{-2}}{kg/m^3\cdot m/s\cdot s^{-2}}}=\sqrt{m^2}=m
$$

(d) Note that $I \propto s_m^2$ so:

$$
10\log\left(\frac{I}{I_0}\right) = 20\log\left(\frac{s_m}{s_{m0}}\right) = \beta
$$

As above, consider dB differences to find s_m ratios:

$$
20\left[\log\left(\frac{s_{m1}}{s_{m2}}\right)\right] = \beta_1 - \beta_2
$$

We are asked to find ratio: final/initial:

$$
\left(\frac{s_{mf}}{s_{mi}}\right) = 10^{(\beta_f - \beta_i)/20} = 10^{-1}
$$

17-50. The Doppler effect formula, Eq. 17-47, and its accompanying rule for choosing ± signs, are discussed in §17-9. Using that notation, we have $v = 343$ m/s, $v_D = v_S = 160000/3600 = 44.4$ m/s, and $f = 500$ Hz. Thus,

$$
f' = (500) \left(\frac{343 - 44.4}{343 - 44.4} \right) = 500 \,\text{Hz} \implies \Delta f = 0 \,.
$$

- 17-61. We use Eq. 17-47 with $f = 500$ Hz and $v = 343$ m/s. We choose signs to produce $f' > f$.
	- (a) The frequency heard in still air is

$$
f' = 500 \left(\frac{343 + 30.5}{343 - 30.5} \right) = 598
$$
 Hz.

(b) In a frame of reference where the air seems still, the velocity of the detector is $30.5 - 30.5 = 0$, and that of the source is $2(30.5)$. Therefore,

$$
f' = 500 \left(\frac{343 + 0}{343 - 2(30.5)} \right) = 608
$$
 Hz.

(c) We again pick a frame of reference where the air seems still. Now, the velocity of the source is $30.5 - 30.5 = 0$, and that of the detector is 2(30.5). Consequently,

$$
f' = 500 \left(\frac{343 + 2(30.5)}{343 - 0} \right) = 589
$$
 Hz.

- 17-63. (a) The half angle θ of the Mach cone is given by $\sin \theta = v/v_s$, where v is the speed of sound and v_s is the speed of the plane. Since $v_s = 1.5v$, $\sin \theta = v/1.5v = 1/1.5$. This means $\theta = 41.8^\circ$.
	- (b) Let h be the altitude of the plane and suppose the Mach cone intersects Earth's surface a distance d behind the plane. The situation is shown on the diagram below, with P indicating the plane and O indicating the observer. The cone angle is related to h and d by $\tan \theta = h/d$, so $d = h/\tan \theta$. The shock wave reaches O in the time the plane takes to fly the distance d: $t = d/v = h/v \tan \theta =$ $(5000 \,\mathrm{m})/1.5(331 \,\mathrm{m/s})\tan 41.8^\circ = 11.3 \,\mathrm{s}.$

